## Efficient I(5 ${}^{2}P_{1/2}$ ) Production by the Reaction of I<sub>2</sub> with NCl<sub>3</sub> Flame

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When gaseous iodine was admitted into the NCl<sub>3</sub> autodecomposition flame, intense I(5  ${}^{2}P_{1/2}$ )  $\rightarrow$  I(5  ${}^{2}P_{3/2}$ ) emission was observed owing to direct pumping of iodine atoms by NCl<sub>3</sub> flame. It is found that the chemical environment of NCl<sub>3</sub> flame is benign for I(5  ${}^{2}P_{3/2}$ ). The mechanisms of I(5  ${}^{2}P_{1/2}$ ) production in the NCl<sub>3</sub>/I<sub>2</sub> reactive gas flow are discussed for the first time. It is believed that Frank–Condon favored *E–E* energy transfer from Cl<sub>2</sub>(B) to I(5  ${}^{2}P_{3/2}$ ) and the nearly resonant *V–E* energy transfer from Cl<sub>2</sub>(X, v = 15) to I(5  ${}^{2}P_{3/2}$ ) respond for the production of excited state iodine atoms.

The first chemically pumped iodine laser oscillating on the I(5  ${}^{2}P_{1/2}$ )–I(5  ${}^{2}P_{3/2}$ ) (hereafter I\* and I) electronic transition was demonstrated by McDermott et al.<sup>1</sup> in 1978. Excited iodine atoms I\* for use in iodine lasers has been produced photolytically<sup>2</sup> and by near resonant energy transfer between ground state iodine atoms and energy carriers, such as singlet oxygen O<sub>2</sub>(a) (chemical oxygen iodine laser, COIL)<sup>3</sup> and NCl(a) (HN<sub>3</sub>-based all gas-phase iodine laser, AGIL).<sup>4</sup> However, as the most successful iodine laser, COIL suffers from the on-demand preparation of basic hydrogen peroxide (BHP) that is unstable to long preservation, from the overall all weight ratio of singlet oxygen generator (SOG), and from water vapor, a relatively strong deactivator of excited iodine atoms.<sup>5a</sup> For HN<sub>3</sub>-based AGILs, the difficulty lies in the highly explosives of gaseous HN<sub>3</sub> when condensed and in a high-temperature combustor to produce F or Cl.<sup>5b</sup>

Recently, McDermott et al.<sup>6a</sup> have observed NCl(a–X) emission spectrum in the  $Cl/H_2/NCl_3$  flow system, and Davis et al.<sup>6b</sup> have observed I\* transition in H/NCl<sub>3</sub>/HI flow system and the resulting significant decrease in the optical absorption of iodine atoms. More recently, McDermott et al.<sup>6c</sup> utilized the NCl<sub>3</sub> autodecomposition reaction as chlorine atom source to remove the necessity of a high-temperature combustor.

However, in spite of many thorough experiments on the NCl<sub>3</sub> system,<sup>6</sup> no laboratory has reported the achievement of laser oscillation based on NCl<sub>3</sub>/H/I reaction schemes. So, further deep studies on NCl<sub>3</sub> reactions to produce excited iodine atoms I<sup>\*</sup> efficiently are needed. In the present study, the direct pumping of iodine atoms by NCl<sub>3</sub> autodecomposition flame was examined and studied.

A bright red flame in the gas-flow reactor due to surfacecatalyzed branched-chain decomposition of NCl<sub>3</sub> was observed downstream the nozzle injector immediately after the gaseous mixture of NCl<sub>3</sub>/He was expanded through the 140- $\mu$ m diameter nozzle injector into the 1–10 Torr flow reactor. The red light is attributed to the B–X emission of Cl<sub>2</sub> molecules and Cl<sub>2</sub>(B) is populated mostly on v = 0 vibrational level. The temperature increment of the reactor is less than 5 °C. This indicates that the wall losses of reactive species can be negligible at experimental conditions and that the energy librated by energetic NCl<sub>3</sub> mole-



Figure 1. IR spectra of  $NCl_3$  combustion flame with and without iodine.



Figure 2. Flashed increase of  $I^*$ -I emission before its extinguishing when  $I_2$  flow is turned off.

cules is mostly distributed on  $Cl_2(B)$  or higher vibrational levels of  $Cl_2(X)$ .

As is shown in Figure 1, intense I<sup>\*</sup>–I emission at 1315 nm is observed immediately after the admission of  $I_2$ /He gas mixture into the NCl<sub>3</sub> flame. It can be seen in Figure 1 that the intense I<sup>\*</sup>–I emission is always accompanied by a little reduction of Cl<sub>2</sub>(B–X) emission intensity.

It is found that when the  $I_2/He$  gas mixture is injected downstream to the NCl<sub>3</sub> flame, no I<sup>\*</sup>–I emission at 1315 nm could be observed. When the  $I_2/He$  gas flow is injected upstream to the NCl<sub>3</sub> flame, the I<sup>\*</sup>–I emission can be observed in the flame and after the flame. These results indicate that excitation of iodine atoms results from Cl<sub>2</sub>(B) energy carrier or some species in NCl<sub>3</sub> flame related with Cl<sub>2</sub>(B). To date, however, there is no reference that discusses the mechanisms of the energy transfer between



**Figure 3.** Rapid production and slow decay of excited iodine atoms  $I^*$ . (a) The time decay of  $I^*$  along the gas-flow reactor; (b) Logarithmic plot for the time decay of  $I^*$ .

NCl<sub>3</sub> flame and iodine atoms.

In the case that  $I_2$  is excessive, an interesting flashedincrease of I<sup>\*</sup>–I emission before its extinguish is observed in our experiments when the  $I_2$  flow is turned off, as is shown in Figure 2. The I<sup>\*</sup> emission is stable under fixed flow rate of NCl<sub>3</sub> and I<sub>2</sub>. However, it was found that the I<sup>\*</sup> emission intensity would always increased sharply and multiply before its final extinguishing. It may be explained by the efficient quenching of I<sup>\*</sup> by I<sub>2</sub>.<sup>7</sup>

Figure 3a shows the rapid production of I<sup>\*</sup> and the I<sup>\*</sup>–I emission intensity decay curve along the flow reactor. Fast production of I<sup>\*</sup> indicates that the energy-transfer process is efficient and the a compact iodine laser is feasible. Figure 3b is a logarithmic plot for the time decay of excited-state iodine atoms. A slow total decay rate,  $k_d = 9.6 \text{ s}^{-1}$ , is obtained by a least-squares fitting. It is comparable with the radiation rate,  $k_r = 7.8 \text{ s}^{-1.7}$  Consequently, when I<sub>2</sub> is almost completely converted into I by reacting with Cl, the quenching decay rate of I<sup>\*</sup>,  $k_q = k_d - k_r$ , is virtually negligible at experimental conditions. These indicate that the chemical environment in the NCl<sub>3</sub> autodecomposition reaction is benign for I<sup>\*</sup> and that a scalable chemical laser system may be feasible.

In order to explain the efficient production of I\*, the following probable mechanisms are proposed.

Iodine atoms can be generated by the rapid reaction of  $I_2$  with Cl atoms in NCl<sub>3</sub> flame. Then, iodine atoms reacting with NCl<sub>3</sub> flame can produce excited iodine atoms  $I^*$ .

$$Cl + I_2 \rightarrow ICl + I \tag{1a}$$

$$Cl + ICl \rightarrow Cl_2 + I$$
 (1b)

One probable mechanism is the energy transfer between energy carriers and iodine atoms. The Frank–Condon factor of  $Cl_2(B, v = 0)-Cl_2(X)^8$  and the energy level diagrams of  $Cl_2(B)$ ,  $Cl_2(X, v)$  and I\* are shown in Figure 4. As is seen in Figure 4, the electronic transition of  $Cl_2(B, v = 0)$  to  $Cl_2(X, v = 15-19)$  is Frank–Condon favored. The electronic transition of  $Cl_2(B, v = 0)$  to  $Cl_2(X, v = 15)$  to  $Cl_2(X, v = 19)$  and vibrational transition of  $Cl_2(X, v = 15)$  to  $Cl_2(X, v = 0)$  is virtually in energy resonance with the electronic transition from I to I\*, and these energy-transfer rates are expected to be extremely efficient. Therefore, it is possible that the high vibrational levels of the X state of  $Cl_2, Cl_2(X, v = 15-19)$  participate in the production of I\*. It includes two pathways, i.e., (i) *E–E* energy transfer:

$$Cl_2(B, v = 0) + I \rightarrow Cl_2(X, v = 19) + I^*$$
(2)  
) *V–E* energy transfer:

and (ii



Figure 4. The most probable mechanism proposed for the production of excited state iodine  $I^*$  in  $NCl_3/I_2$  reactive flow system.

$$Cl_2(B, 0) \to Cl_2(X, v = 15) + hv(1007 \text{ nm})$$
 (3a)

$$Cl_2(X, v = 15) + I \rightarrow Cl_2(X, v = 0) + I^*$$
 (3b)

As aforementioned, the flow tube studies show that the intense I<sup>\*</sup> emission could be produced only in the presence of  $Cl_2(B)$ . It also can be seen in Figure 1 that the intense I<sup>\*</sup>–I emission is always accompanied by a little reduction of  $Cl_2(B-X)$  emission intensity. These suggest that  $Cl_2(B)$  or some excited state species related with  $Cl_2(B)$  must participate in the efficient production of I<sup>\*</sup>.

Another probable mechanism may be due to  $Cl_3(AL)$ . Linear type  $Cl_3$  intermediate has a strongly bound state AL, located about 8900 cm<sup>-1</sup> above the ground state XL, which is in nearly resonance with I\*–I transition.<sup>9</sup> The  $Cl_3(AL)$  intermediate may be produced in NCl<sub>3</sub> flame by the recombination of Cl and  $Cl_2(B)$ . More mechanisms can be seen in detail in Supporting Information.<sup>10</sup>

In summary, intense I\*–I emission could be observed by the reaction of I<sub>2</sub> with NCl<sub>3</sub> flame. I\* was produced only in the presence of Cl<sub>2</sub>(B). The chemical environment of the NCl<sub>3</sub> autode-composition flame is benign for I\*. The NCl<sub>3</sub>/I<sub>2</sub>/He system can be a scale efficient iodine laser system. Finally, the most probable mechanisms for I\* production are proposed for the first time. They include nearly resonant energy-transfer processes from Cl<sub>2</sub>(B) and Cl<sub>2</sub>(X, v = 15) to I(5 <sup>2</sup>P<sub>3/2</sub>).

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- 10 Supporting Information is available electronically on the CSJ-Journal Web site, http://www.csj.jp/journals/chem-lett/index.html.